

THE $U_A(1)$ ANOMALY AND THE η' MASS IN COULOMB GAUGE QCD

DIPLOMA THESIS

Gernot Lassnig

Institute for Physics, SIC!QFT Group
University of Graz

Supervisor *Reinhard Alkofer & Diana Nicmorus*

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CONTENTS

1 MOTIVATION

- QCD Symmetries
- $U_A(1)$ Anomaly

2 η' MASS GENERATION

- Coulomb Gauge

3 DYSON-SCHWINGER AND BETHE-SALPETER EQUATIONS

4 CURRENT AND FUTURE WORK



FUNDAMENTALS OF QCD

QCD LAGRANGIAN

$$\mathcal{L}_{\text{QCD}} = \sum_{f=1}^{N_f} \bar{\psi}^f (i\gamma_\mu D^\mu - m_f) \psi_f - \frac{1}{2} \text{tr}(F_{\mu\nu} F^{\mu\nu}) + \mathcal{L}_{\text{ghost}} + \mathcal{L}_{g.f}$$

Left and Right handed fields

$$\psi_R^f = \frac{1}{2}(1 + \gamma_5)\psi^f$$

$$\Rightarrow \bar{\psi}^f i\gamma_\mu D^\mu \psi^f = \bar{\psi}_R^f i\gamma_\mu D^\mu \psi_R^f + \bar{\psi}_L^f i\gamma_\mu D^\mu \psi_L^f$$

$$\psi_L^f = \frac{1}{2}(1 - \gamma_5)\psi^f$$



FUNDAMENTALS OF QCD

SYMMETRY GROUP OF MASSLESS LAGRANGIAN

$$U_L(N_f) \times U_R(N_F) \cong U_{L+R}(1) \times U_{L-R}(1) \times SU_L(N_F) \times SU_R(N_F)$$

$$U_{L+R}(1)$$

Conservation of baryon number

$$U_{L-R}(1)$$

Axial symmetry

$$SU_L(N_f) \times SU_R(N_f)$$

Chiral symmetry, spontaneously broken

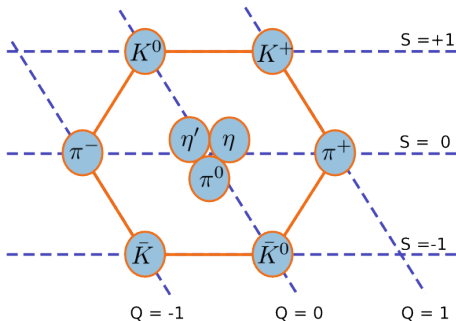
Hadronic vacuum only invariant under $SU(N_F)$

Goldstone theorem $\Rightarrow N_f^2 - 1$ **massless** Goldstone Bosons



$U_A(1)$ ANOMALY (SEE RICHARD WILLIAMS TALK)

- $U_A(1)$ not a symmetry of nature, so it's **broken**
- Upper limit for Goldstone Boson is $\sqrt{3}m_\pi$ ¹
- No particle known



¹S.Weinberg

$U_A(1)$ ANOMALY

- $U_A(1)$ Symmetry is “anomalously” broken
- Feynman graphs containing quark triangles create anomaly (ABJ)
- **No** conservation of axial vector current

AXIAL CURRENT

$$j_\mu^5 = \bar{\psi}_f \gamma_\mu \gamma^5 \psi_f$$

$$\partial^\mu J_\mu^5 = \sum_f 2im_f \bar{q}_f \gamma^5 q_f + N_f \frac{g^2}{16\pi^2} F^{a\mu\nu} \tilde{F}_{\mu\nu}^a$$

Not conserved but . . .



$U_A(1)$ ANOMALY

$N_f \frac{g^2}{16\pi^2} F^{a\mu\nu} \tilde{F}_{\mu\nu}^a$ can be written as 4 divergence
 \Rightarrow new **conserved** current $\tilde{J}_\mu^5 = J_\mu^5 - \frac{N_f g^2}{2\pi^2} K_\mu$

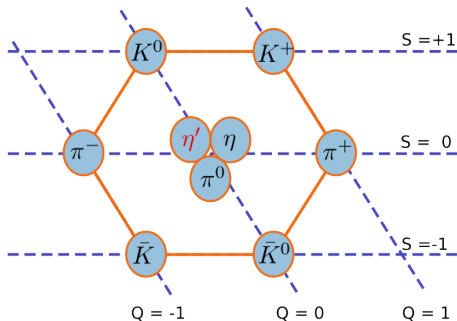
CONSERVED CHARGE

$$\frac{d\tilde{Q}_5}{dt} = \int d^3x \partial_0 \tilde{J}_0^5 = 0 \quad (1)$$

Looks like $U_A(1)$ problem persist

- General: K_μ not gauge invariant, so \tilde{J}_μ^5 not gauge invariant
- t'Hofft: Equation (1) not correct in presence of instantons



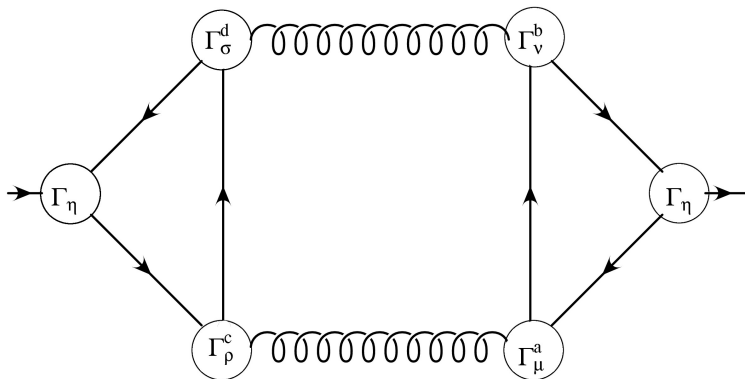
η' MASS GENERATION

η' is $SU(3)$ singlet state \Rightarrow mixes with 2 gluons
 Diamond diagram can contribute to η' mass ¹

¹Diploma Thesis, Almut Mecke, 1997 & R. Williams, C.S. Fischer, R. Alkofer to be published

LOWEST ORDER DIAGRAM

Diamond diagram



WHY IS COULOMB GAUGE INTERESTING

Working in Coulomb gauge ($\vec{\nabla} \cdot \vec{A} = 0$)

- advantages

- ▶ Coulomb gauge is physical gauge
- ▶ All degrees of freedom are physical
- ▶ No confinement without Coulomb Confinement ¹
- ▶ Working directly in Minkowski space

- disadvantages

- ▶ No proof of renormalisability
- ▶ Complicated calculations

¹D.Zwanziger, Phys. Rev. Lett. 90, 102001 (2003)



DYSON-SCHWINGER AND BETHE-SALPETER EQUATIONS

Quark self-energy due to gluons

QCD GAP EQUATION

$$iS^{-1}(p) = \not{p} - m - \Sigma(p) \quad (2)$$

A $q\bar{q}$ bound state described by the Bethe-Salpeter equation

$$\Gamma(P, q) = \int d^4k K(q, k, P) S(k_+) \Gamma(P, K) S(k_-)$$

$\Gamma(P, q)$ Bethe-Salpeter Amplitude

Infinite coupled system \Rightarrow truncation needed



DYSON-SCHWINGER AND BETHE-SALPETER EQUATIONS

$$\Sigma(p) = C_f 6\pi \int \frac{d^4 q}{(2\pi)^4} V_c(\vec{k}) \gamma_0 S(q) \gamma_0 \quad (3)$$

Ansatz

$$S^{-1}(p) = -i [\gamma_0 p_0 - \vec{\gamma} \cdot \vec{p} C(p) - B(p)] \quad (4)$$

leads to two coupled integral equations

$$B(p) = m + \frac{1}{2\pi^2} \int d^3 q V_c(k) \frac{M(q)}{\tilde{\omega}(q)}$$

$$C(p) = 1 + \frac{1}{2\pi^2} \int d^3 q V_c(k) \hat{p} \cdot \hat{q} \frac{q}{p \tilde{\omega}(q)}$$

$$M(p) := \frac{B(p)}{C(p)}$$

$$\tilde{\omega}(q) = \sqrt{M(q)^2 + p^2}$$

$$C_f = 4/3$$



QUARK MASS FUNCTION

Highly singular non-linear coupled integral equation

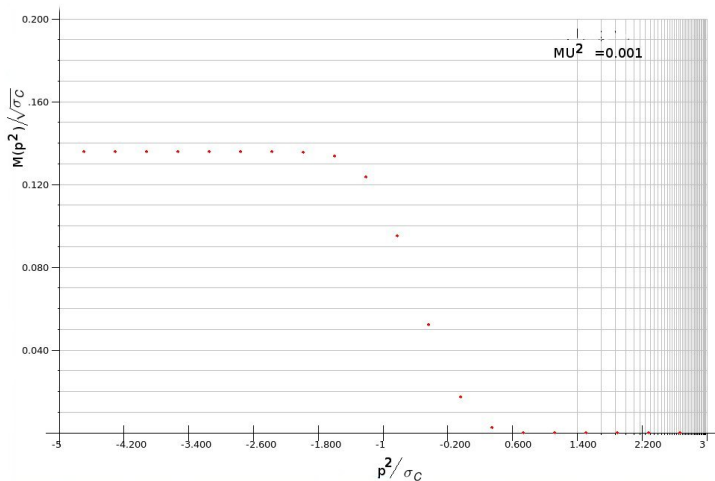
- Iterating is the only known solution
- Gauß-Kronrod Integration system is used
- Coulomb Gluon Part $V_c = \frac{\sigma_c}{(k^2)^2}$
 \Rightarrow Angular integrations can be done analytically

$$B(p) = m + \frac{1}{2\pi^2} \int d^3q V_c(k) \frac{M(q)}{\tilde{\omega}(q)}$$

$$C(p) = 1 + \frac{1}{2\pi^2} \int d^3q V_c(k) \hat{p} \cdot \hat{q} \frac{q}{p\tilde{\omega}(q)}$$

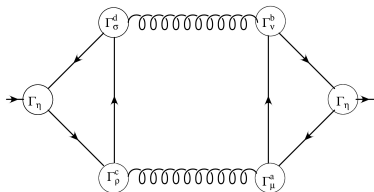


SOLUTION OF THE COUPLED INTEGRAL EQUATION



CURRENT WORK

Power counting



Chiral limit \Rightarrow BSA results in $\gamma_5 B(q^2)/f_\pi$

Quark propagator known in terms of $B(q)$ and $C(q)$

Working directly in Minkowski space



OUTLOOK

- Test of Quark-Gluon-Vertex: IR singularities needed?
- Using full Bethe-Salpeter Amplitude
- Higher order Diagrams?
- ... lots of stuff to do

